

## SHORT COMMUNICATIONS

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*Acta Cryst.* (1983). **A39**, 422

**Conditions for direct structure imaging in silicon carbide polytypes. Erratum.** By DAVID J. SMITH and M. A. O'KEEFE, *High Resolution Electron Microscope, University of Cambridge, Free School Lane, Cambridge CB2 3RQ, England*

(Received 11 February 1983)

**Abstract**

As a result of a printer's error, Figs. 9(a) and 9(b) in Smith & O'Keefe [*Acta Cryst.* (1983), **A39**, 139–148] have been transposed, so that (a) is labelled (b) and (b) is labelled (a).

All information is given in the *Abstract*.

*Acta Cryst.* (1983). **A39**, 422–424

**A cusp-constrained scattering factor for bonded hydrogen atoms.** By ROB J. VAN DER WAL and ROBERT F. STEWART,\* *Department of Chemical Physics, University of Groningen, Nijenborgh 16, 9747 AG Groningen, The Netherlands*

(Received 19 October 1982; accepted 7 December 1982)

**Abstract**

A simple one-parameter form factor that satisfies a cusp condition has been developed. A one-parameter form factor, based on a 1s-type density function, is more tractable in a least-squares analysis than the corresponding cusp-constrained form factor. Application of both form factors for bonded hydrogen atoms to X-ray diffraction data of sucrose revealed a more contracted density for H bonded to C compared to H bonded to O.

At a NATO-sponsored international school on electron and magnetization densities in molecules and crystals (Becker, 1980) held in Arles, France, in August 1978, a number of exercises were distributed by the participants. One exercise, proposed by Professor Hirshfeld, was to develop an atomic form factor which satisfies the electron–nuclear cusp condition. The present communication is an approximate solution to the problem for a bonded hydrogen atom.

The electron–nuclear cusp condition,

$$\lim_{r \rightarrow 0} [(\partial \bar{\rho} / \partial \mathbf{r}_j) / \bar{\rho}] = -2Z_j, \quad (1)$$

is a consequence of satisfying the many-electron *non-relativistic* Schrödinger equation (Kato, 1957; Steiner, 1963; Pack & Brown, 1966). In (1)  $\bar{\rho}$  is the one-electron density spherically averaged about the  $j$ th nucleus with atomic number  $Z_j$ . The  $r_j$  is a scalar radial length in atomic units (1 a.u. = 0.529177 Å). One possible form for  $\rho$  is

$$\rho(\mathbf{r}) = A_0 e^{-\gamma r} \sum_{n=0}^N (A_n/A_0) r^n / 4\pi. \quad (2)$$

For (2) to satisfy (1),

$$A_1/A_0 = \gamma - 2Z. \quad (3)$$

For  $\rho(\mathbf{r})$  normalized to unity,

$$A_0 = 1 / \left[ (2/\gamma)^3 (1 - 3Z/2\gamma) + \sum_{n=2}^N (A_n/A_0) (n+2)! / \gamma^{n+3} \right]. \quad (4)$$

To maintain simplicity, we set all  $A_n/A_0$  equal to zero for  $n \geq 2$ . This gives a simple one-parameter function which normalizes to one and satisfies (1):

$$\rho(\gamma, \mathbf{r}) = (\gamma/2)^3 (1 - 3Z/2\gamma)^{-1} e^{-\gamma r} [1 + (\gamma - 2Z) r] / 4\pi. \quad (5)$$

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